



# Environment expansion for Matrix Product States

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NATIONAL TSING HUA UNIVERSITY



國家科學及技術委員會  
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  - Accelerating subspace expansion (post-expansion)
- 4 Benchmark example

# Introduction: DMRG

General many-body quantum state as a tensor product of  $d$ -dimensional local Hilbert spaces:

$$|\Psi\rangle = \underbrace{\hspace{10em}}_{c^{s_1 \dots s_N}}, \quad \mathcal{O}(\exp N) \text{ DOFs.}$$

Compress as a **matrix product state**:

$$|\Psi\rangle = \begin{array}{ccccccc} & s_1 & s_2 & s_3 & \dots & s_N & \\ & \circ & \circ & \circ & \dots & \circ & \\ & | & | & | & & | & \\ & A_1^{s_1} & A_2^{s_2} & A_3^{s_3} & \dots & A_N^{s_N} & \end{array}, \quad \mathcal{O}(ND^2) \text{ DOFs.}$$

Good at representing locally entangled states.  
Accuracy controlled by  $D$ .



# Density Matrix Renormalization Group

See Schollwöck, Annals of Physics 326, 96 (2011) for a review  
doi:10.1016/j.aop.2010.09.012

The 'classic' algorithm is 2-site DMRG: update two sites of the tensor network at once.

One DMRG step: 2 sites,  $Dd \times dD$  dimensional tensor

$$|\Psi\rangle = \text{---} \text{---} \text{---} \quad H |\Psi\rangle = \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array}$$

Cost of the tensor network contraction is  $\mathcal{O}(wd^2D^3)$

SVD to convert  $\psi$  back into MPS form:

$$|\Psi\rangle = \text{---} \text{---} \text{---} = \text{---} \text{---} \text{---}$$

Cost of the SVD is  $\mathcal{O}(d^3D^3)$

# Why not update a single site?

- Much faster –  $\mathcal{O}(wdD^3)$  for the matrix-vector multiply

$$|\Psi\rangle = \text{---} \bigcirc \text{---} \quad H |\Psi\rangle = \left( \begin{array}{c} \text{---} \\ \text{---} \end{array} \right) \bigcirc \left( \begin{array}{c} \text{---} \\ \text{---} \end{array} \right)$$

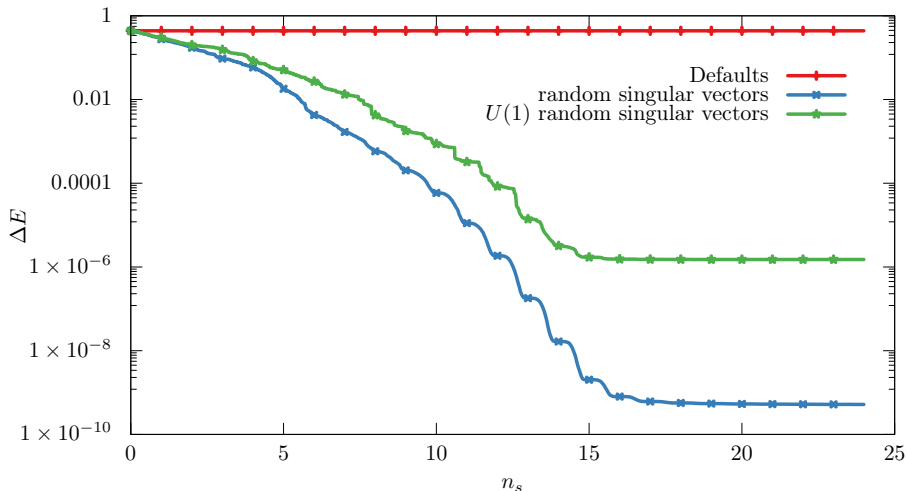
- SVD is  $\mathcal{O}(dD^3)$

$$\text{---} \bigcirc \text{---} = \text{---} \bigcirc \text{---} \square \text{---}$$

- *Linear* scaling in the local Hilbert space dimension  $d$
- Problem: how to increase the bond dimension? (random states?)

# Single-site convergence

Heisenberg spin chain example – Néel state to  $D = 100$

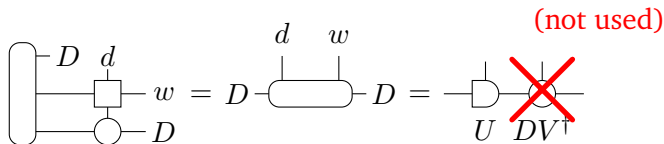


With  $U(1)$  symmetry, the dimension in each symmetry sector is *frozen*

# Single-site subspace expansion (3S)

C. Hubig, IPM, U. Schollwöck, and F. A. Wolf, Phys. Rev. B 91, 155115 (2015)  
doi:10.1103/PhysRevB.91.155115

- Incorporate *new degrees of freedom* during the truncation using the environment Hamiltonian



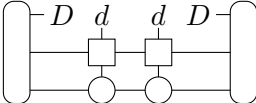
Cost of contraction is  $\mathcal{O}(wdD^3)$ . Cost of SVD is  $\mathcal{O}(wd^2D^3)$ .

Earlier method from Steven R White that uses an eigenvalue decomposition,  $\mathcal{O}(d^3D^3)$  [Phys. Rev. B 72, 180403 (2005)]



## 3S – why does it work?

Easier to understand as a density matrix

$$\rho' = \sum_w E_w \rho E_w^\dagger = X X^\dagger =$$


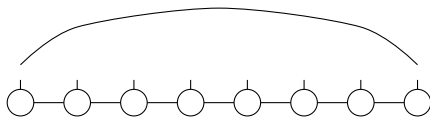
Perturb the density matrix using the block operators

- It is important to use appropriate weights
- One block operator is the identity
- The weight of the remaining components is the control parameter  $\alpha$
- Weight each component by the norm  $F_w$  matrix
- Omit the block Hamiltonian itself – don't need it
- The remaining set  $\{E_w\}$  contains all interaction terms that are needed at the current site
  - They might not appear in the projected Hamiltonian itself!

# Adding relevant degrees of freedom

- Long-range interactions

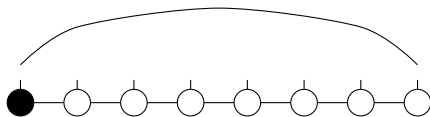
Example: periodic boundary conditions



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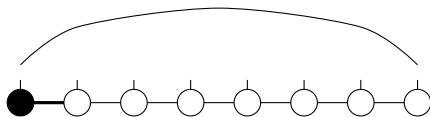


Matrix elements are zero at the active site

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Example: periodic boundary conditions

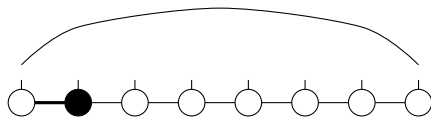


Relevant basis states at site 1 added via the mixing term

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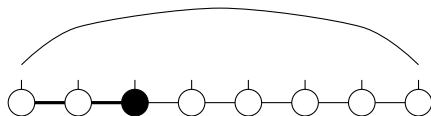


Relevant basis states are kept throughout the sweep

# Adding relevant degrees of freedom

- Long-range interactions

Example: periodic boundary conditions

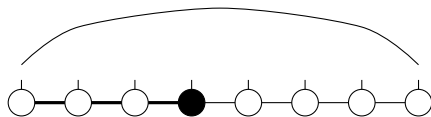


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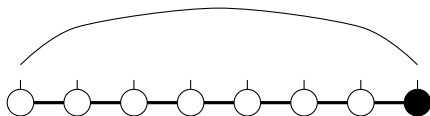


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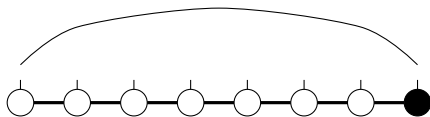
Now matrix elements are non-zero



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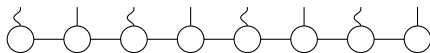


Now matrix elements are non-zero

- Inhomogeneous degrees of freedom

Example: fermions coupled to bosonic degrees of freedom

$$H = -t \sum_{\langle i,j \rangle, \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.}) + U \sum_i n_{i\uparrow} n_{i\downarrow} + \omega \sum_i b_i^\dagger b_i + g \sum_i (n_{i\uparrow} + n_{i\downarrow} - 1) (b_i^\dagger + b_i)$$



2-site update cannot introduce new quantum number sectors

# Low rank factorization: randomized SVD (RSVD)

N. Halko, P.-G. Martinsson, and J. A. Tropp, SIAM Rev. 53, 217 (2011)

doi:10.1137/090771806

- We only want a fraction  $\sim 1/d$  of the singular vectors.
- Traditional algorithms calculate *all* singular values and truncate.

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## Range finding algorithm

Find (approximate) dominant  $k$  left singular values of a  $m \times n$  matrix  $M$

Construct  $n \times k$  Gaussian random matrix  $\Omega$

QR decomposition:

$$QR = M\Omega$$

$Q$  is an  $n \times k$  matrix containing the dominant  $k$  (approximate) left singular vectors

- Cost is dominated by the matrix multiply  $M\Omega$

# Randomized SVD (RSVD)

Improving accuracy: *oversampling*

## Randomized SVD

$p$  is the *oversampling parameter* ( $p = 10$  is typical)

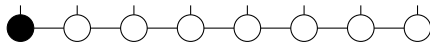
- 1 Construct  $\Omega$ ,  $n \times (k + p)$  Gaussian random matrix
- 2  $QR$  decomposition:  $QR = M\Omega$   $Q$  is  $n \times (k + p)$
- 3 Singular value decomposition  $UDV^\dagger = Q^\dagger M$
- 4 Keep the  $k$  largest singular values  $U$  is  $(k + p) \times k$
- 5  $QU$  is a good approximation to the dominant  $k$  singular values

Oversampling does not increase cost – dominated by multiplication  $M\Omega$   
Very good theoretical error bounds with modest oversampling  
 $p = 10$  works OK.

# Pre-expansion versus post-expansion

We refer to two-site DMRG and related methods as *pre-expansion*

- Bond expansion *before* the optimization step
- Incorporate degrees of freedom from the environment
- These degrees of freedom are ‘thrown away’ afterwards

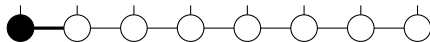


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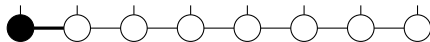


*Expand the bond dimension*

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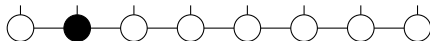


*Optimize the site, then truncate*

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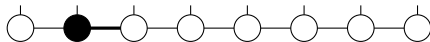
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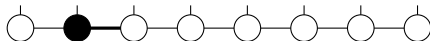


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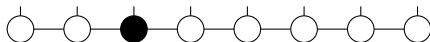


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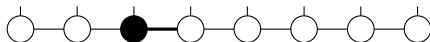


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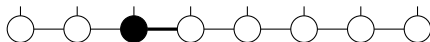


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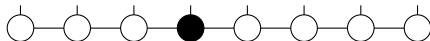


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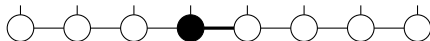


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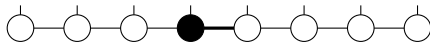


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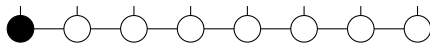
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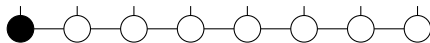


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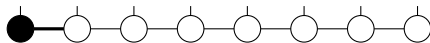


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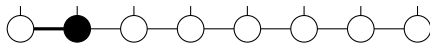


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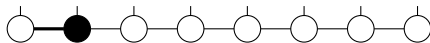


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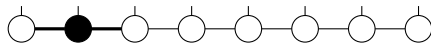


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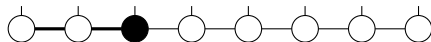


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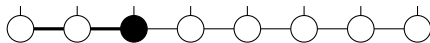


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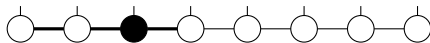
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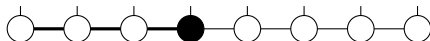


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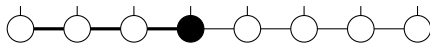


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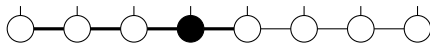


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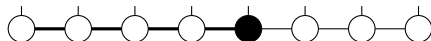


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*Next site*

# Accelerating 2-site DMRG (pre-expansion)

- Kohn, Tschirsich, Keck, Plenio, Tamascelli, Mongangero, Phys. Rev. E 97, 013301 (2018)  
doi:10.1103/PhysRevE.97.013301  
2-site DMRG using RSVD,  $\mathcal{O}(d^2 D^3)$  no longer a bottleneck for large  $d$
- Andreas Gleis, Jheng-Wei Li, and Jan von Delft, Phys. Rev. Lett. 130, 246402 (2023)  
doi:10.1103/PhysRevLett.130.246402  
2-site DMRG with single site cost, overall  $\mathcal{O}(dwD^3)$ , but a total of 5 SVD's

Basic idea:

- One iteration of 2-site DMRG, enlarge bond dimension by  $k$
- Single-site DMRG in this  $D + k$  dimensional basis

Overall cost for the optimization step is  $\mathcal{O}(dw(D + k)^3)$

## Accelerated 2-site DMRG

Simpler and faster contraction versus the method from Gleis *et al*

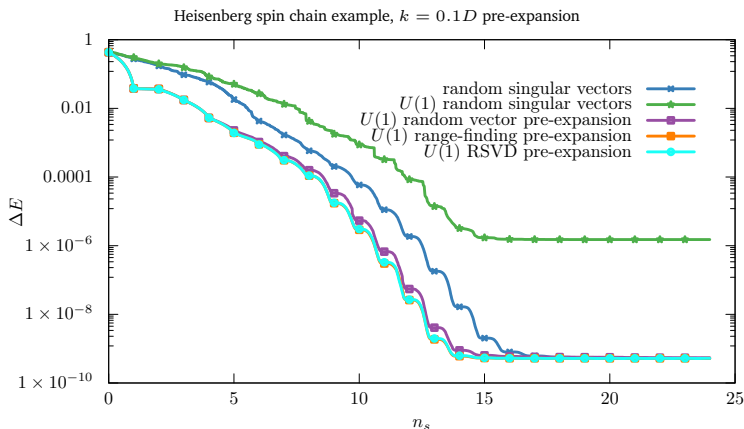
$$M\Omega = \left( \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right) \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \left( \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right)$$

Contraction cost is  $\mathcal{O}(dwkD^2)$ . No  $\mathcal{O}(D^3)$  operations.

- 1 Construct  $dD \times k$  Gaussian random matrix  $\Omega$
- 2 Orthogonalize  $\Omega$  against the states already in the environment basis
- 3 Insert this into the 2-site matrix-vector multiply
- 4  $QR$  decomposition  $QR = M\Omega$
- 5  $Q$  is  $dD \times k$  matrix of the expansion vectors
- 6 Augment the environment site and Hamiltonian

# Pre-expansion variants

- Random – choose random expansion vectors
  - Same distribution of quantum number sectors as the existing states
  - Ensure at least one state in each available sector
- Range-finding – Use  $Q$  from  $QR = M\Omega$  as the expansion vectors
- RSVD – Oversample the range-finding algorithm and SVD of  $Q^\dagger M$





# Post-expansion

Apply the RSVD ideas to the 3S algorithm

- Partition the states into two classes
  - $D$  kept states
  - $k$  additional *expansion vectors*
  - the MPS has  $D + k$  states total
- Orthogonalize the expansion vectors against the kept states
  - only need  $k$  singular vectors of the big matrix, rather than  $D$  (or  $D + k$ )

## Accelerated 3S mixing

$$M\Omega = \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \left( \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \right) \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \Omega$$

Contraction cost is  $\mathcal{O}(dwkD^2)$

Weight the components  $\Omega$  from the norm of the  $F$  matrices

# Post-expansion

$$M\Omega = \left( \begin{array}{c} \text{---} \\ E \\ \text{---} \end{array} \right) \begin{array}{c} \text{---} \\ \square \\ \text{---} \\ \text{---} \\ \square \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \Omega \\ \text{---} \end{array} \quad k$$

Physical picture: we don't know (or don't trust) the  $F$  matrix elements – replace them with random numbers!

# Post-expansion

$$M\Omega = \left( \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} E \right) \begin{array}{c} \text{---} \\ | \\ \text{---} \\ | \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ | \\ \text{---} \\ | \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ | \\ \text{---} \\ | \\ \text{---} \end{array} \Omega \end{array}$$

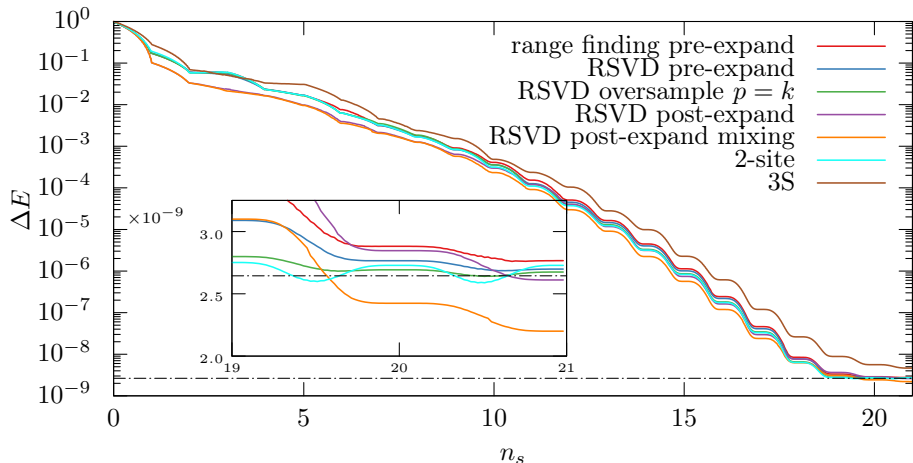
Physical picture: we don't know (or don't trust) the  $F$  matrix elements – replace them with random numbers!

## Post-expansion variants

- Random – choose random expansion vectors
  - Usually not effective – similar to single-site with no singular value cutoff
- Range-finding – Use  $Q$  from  $QR = M\Omega$  as the expansion vectors
  - No control over quantum number sectors
- RSVD – Oversample the range-finding algorithm and SVD of  $Q^\dagger M$
- Mixing – merge the expansion vectors with a mixing factor
  - Set the mixing factor from the discarded weight of the truncation from  $D + k$  to  $D$  – avoids the biggest problem of the old 3S algorithm

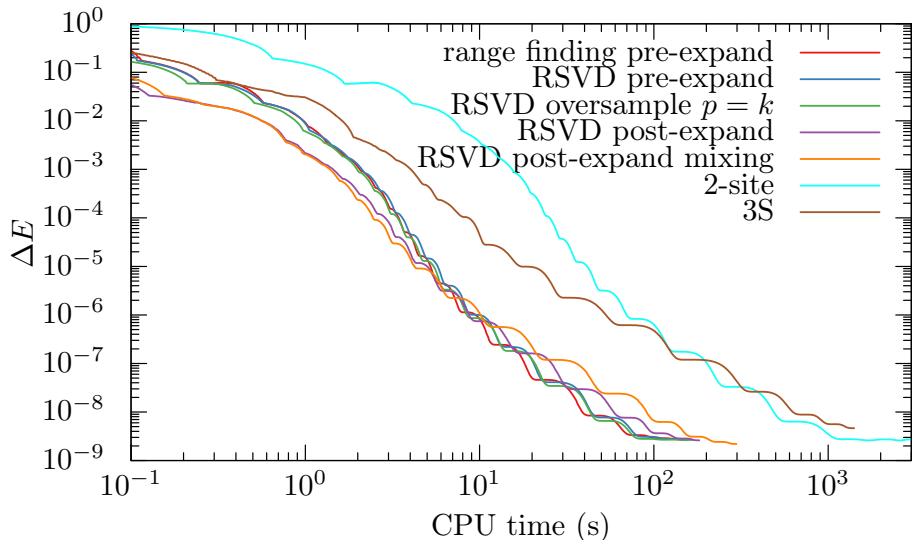
# Benchmark – Hubbard-Hoitein example

$$H = -t \sum_{\langle i,j \rangle, \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + \text{H.c.}) + U \sum_i n_{i\uparrow} n_{i\downarrow} + \omega \sum_i b_i^\dagger b_i + g \sum_i (n_{i\uparrow} + n_{i\downarrow} - 1) (b_i^\dagger + b_i)$$



Convergence is better than 3S

# Hubbard-Hoitein example – CPU time



# Conclusion

Two bond expansion methods for DMRG

*Pre-expansion:*

- Similar convergence as 2-site DMRG
- Same performance as single-site DMRG; linear in  $d$
- Bond expansion is *fast* – effectively zero cost

*Post-expansion:*

- Similar to 3S, good for inhomogeneous or long-range interactions
- Converges better than 3S
- Bond expansion is asymptotically faster than 3S (zero cost)
- More details at arXiv:2403.00562
- Long paper in preparation

Code available: <https://github.com/mptoolkit>

Documentation: <https://mptoolkit.qusim.net>